Correcting the extended-source calibration for the Herschel-SPIRE Fourier-transform spectrometer

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ABSTRACT
We describe an update to the Herschel-Spectral and Photometric Imaging Receiver (SPIRE) Fourier-transform spectrometer (FTS) calibration for extended sources, which incorporates a correction for the frequency-dependent far-field feedhorn efficiency, $\eta_{ff}$. This significant correction affects all FTS extended-source calibrated spectra in sparse or mapping mode, regardless of the spectral resolution. Line fluxes and continuum levels are underestimated by factors of 1.3–2 in the spectrometer long wavelength band (447–1018 GHz; 671–294 $\mu$m) and 1.4–1.5 in the spectrometer short wavelength band (944–1568 GHz; 318–191 $\mu$m). The correction was implemented in the FTS pipeline version 14.1 and has also been described in the SPIRE Handbook since 2017 February. Studies based on extended-source calibrated spectra produced prior to this pipeline version should be critically reconsidered using the current products available in the Herschel Science Archive. Once the extended-source calibrated spectra are corrected for $\eta_{ff}$, the synthetic photometry and the broad-band intensities from SPIRE photometer maps agree within 2–4 per cent – similar levels to the comparison of point-source calibrated spectra and photometry from point-source calibrated maps. The two calibration schemes for the FTS are now self-consistent: the conversion between the corrected extended-source and point-source calibrated spectra can be achieved with the beam solid angle and a gain correction that accounts for the diffraction loss.

Key words: instrumentation: spectrographs – space vehicles: instruments – techniques: spectroscopic.

1 INTRODUCTION
The calibration of an instrument consists of two tasks: (i) removing all instrument signatures from the data and (ii) converting the...
products to physical units using a suitable calibration schema. For
the first task, a good knowledge of the instrument and its response
to different conditions (e.g. observing mode, internal and external
thermal and radiation environments, the solar aspect angle, etc.)
is required. For the second task, a calibration source of assumed
flux or temperature is used to convert the measured signal to physi-

cally meaningful units. The atmosphere blocks most far-infrared
radiation from reaching the ground, therefore the calibration of
far-infrared space borne instrumentation requires a bootstrapping
approach based on previous observations and theoretical models
of candidate sources, typically planets or asteroids.

An imaging Fourier-transform spectrometer (FTS) is part of
the Spectral and Photometric Imaging Receiver (SPIRE; Griffin
et al. 2010) on board the Herschel Space Observatory (Pilbratt
et al. 2010). SPIRE is one of the most rigorously calibrated far-
infrared space instruments to date. It underwent five ground-based
test campaigns and regular calibration observations during the
nearly 4 yr of in-flight operations of Herschel. The stable space
environment at the second Lagrange point and the flawless opera-
tion of the instrument resulted in unprecedented accuracy both in
terms of the telescope and instrument response. A detailed descrip-
tion of the FTS instrument and its calibration scheme is provided in
Swinyard et al. (2010), with an update in Swinyard et al. (2014).

There are no prior systematic studies of the extended-source cal-
ibration for the FTS. Extended-source calibrated maps from the
SPIRE photometer, corrected to the absolute zero level derived
via cross-calibration with Planck-HFI (Bertincourt et al. 2016),
became available during the post-operations phase of Herschel.
These maps allowed for a detailed comparison between photom-
etry and spectroscopy of extended sources. Initial checks showed
significant and systematic differences at levels of 40–60 per cent
across the three photometer bands. Some authors also reported
discrepancies (Kamenetzky et al. 2014; Köhler et al. 2014) and
implemented corrections in order to match the spectra with the
photometry. Others proceeded by starting from the point-source
 calibration and correcting for the source size (e.g. Kamenetzky
et al. 2015; Wu et al. 2015; Makiwa et al. 2016; Morris et al. 2017;
Schirm et al. 2017).

The reported differences with the photometer did not initially
draw our attention, because the comparison is intricate and depends
on the assumptions made. As shown in Wu et al. (2013), the cou-
pling of sources that are neither point-like nor fully extended (i.e.
semi-extended) require good knowledge of the FTS beam and its
side-lobes, as well as good knowledge of the source brightness
distribution. Even extended sources with significant sub-structure
couple in a complicated way with the multimodal and non-Gaussian
beam (Makiwa et al. 2013). Moreover, the source size would imply
colour-correcting the photometry (see Valtchanov 2017. The SPIRE
Handbook, section 5.8; H17 from now on). Hence, both sides of
the comparison need their proper corrections.

In this study, we have tried to alleviate some of the uncertainties
by carefully selecting truly extended sources for cross-comparison
with broad-band intensities from the SPIRE photometer extended-
source calibrated maps. The results of this analysis show a signif-
cicant correction is needed in order to match the extended-source
calibrated spectra with the photometry. This paper introduces the
methods used to derive the necessary corrections, demonstrates the
self-consistency between FTS point and extended-source calibrated
spectra, and demonstrates a good agreement with broad-band pho-

tometry from the SPIRE photometer.

Herschel’s two other instruments, the Heterodyne Instrument
for the Far Infrared (HIFI; de Graauw et al. 2010) and the Photodetector

Array Camera and Spectrometer (PACS, Poglitsch et al. 2010),
share some spectral overlap with the SPIRE FTS. Analysis of a
sample of calibration targets has shown an overall agreement of
±20 per cent between the SPIRE FTS and HIFI, and discrepancies
up to a factor of 1.5–2 for comparisons with PACS (Puga et al.,
in preparation). Noting that the instantaneous bandwidth of HIFI
(2.4 or 4 GHz depending on observing mode and band) is only
marginally wider than the instrumental line shape of the SPIRE
FTS (1.2 GHz), the overall agreement between HIFI and the SPIRE
FTS is acceptable. The spectral overlap between the SPIRE FTS
and the PACS spectrometer falls in 194–210 μm, which is an area
affected by a PACS spectral leak (see Vandenbussche et al. 2016).
Although we have performed a comparison between instruments
for a sample of extended sources, some results were inconclusive
and we have not included this work in this paper.

The structure of the paper is as follows. In Section 2 we briefly
outline the extended-source calibration scheme. In Section 3 we
compare FTS results with photometry from SPIRE maps using a
selection of spatially extended sources and derive a correction that
matches the known far-field feedhorn efficiency. In Section 4 we
link the two FTS calibration schemes (i.e. the point-source and
the corrected extended-source schemes) using the beam solid angle
and a correction for diffraction loss. Some guidelines on using
the corrected spectra are presented in Section 5. In Section 6 we
outline the significance of the correction and the impact on deriving
physical conditions if the uncorrected spectra are used. In Section 7
we present the conclusions.

As much as possible we follow the notations used in the SPIRE
Handbook (H17). Throughout the paper we interchangeably use
intensity and surface brightness as equivalent terms, in units of
either (MJy sr⁻¹) or (W m⁻² Hz⁻¹ sr⁻¹).

2 Telescope Model Based
Extended-Source Calibration

In the following, we briefly outline the main points in the FTS
calibration scheme, which is presented in greater detail in Swinyard
et al. (2014).

As there is no established absolute calibration source for extended
emission in the far-infrared and sub-mm bands, the Herschel tele-
scope itself is used as a primary calibrator for the FTS. The usual
sources used from ground, such as the Moon and the big planets
(e.g. Wilson, Rohlfs & Hüttemeister 2013), are either too close to
the Sun/Earth or too bright for the instrument.

The SPIRE FTS simultaneously observes two very broad over-
lapping spectral bands. The signals are recorded with two arrays
of hexagonally close-packed, feedhorn-coupled, bolometer detectors:
the spectrometer short wavelength (SSW) array with 37 bolome-
ters, covering 191–318 μm (1568–944 GHz) and the spectrometer
long wavelength (SLW) array with 19 bolometers, covering 294–
671 μm (1018–447 GHz). The bolometers operate at a temperature
of ∼300 mK, which is achieved with a special ³He sorption cooler
(see H17 for more details).

Within the FTS, the radiation from the combination of the astro-
nomical source, the telescope, and the instrument² is split into two

¹ 1 MJy sr⁻¹ = 10⁻²⁰ W m⁻² Hz⁻¹ sr⁻¹.
² The instrument contribution enters in the total radiation because of the Mach–Zehnder configuration of the FTS, where a second input port views
an internal blackbody source (see H17 for more details).
beams. A moving mirror introduces an optical path difference between the two beams. The recombination of the beams produces an interferogram on each of the individual feedhorn-coupled bolometers. Hence, the recorded signal \( V_{\text{obs}} \) after Fourier transforming the interferograms, can be expressed as

\[
V_{\text{obs}} [\text{V Hz}^{-1}] = R_{S} I_{S} + R_{\text{tel}} M_{\text{tel}} + R_{\text{inst}} M_{\text{inst}},
\]

(1)

where \( I_{S} \) is the source intensity, \( M_{\text{tel}} \) and \( M_{\text{inst}} \) are the intensities corresponding to the telescope and the instrument emission models, \( R_{S} \), \( R_{\text{tel}} \), and \( R_{\text{inst}} \) are the relative spectral response functions (RSRFs) of the system for the source, the telescope, and the instrument, respectively. We assume the instrument and telescope emissions to be fully extended in the beam, and well represented by blackbody functions and \( R_{S} = R_{\text{tel}} \). The units of \( I_{S}, M_{\text{tel}}, \) and \( M_{\text{inst}} \) are [W m\(^{-2}\) Hz\(^{-1}\) sr\(^{-1}\)], therefore the RSRF are in units of [V Hz\(^{-1}\)/(W m\(^{-2}\) Hz\(^{-1}\) sr\(^{-1}\)].

The instrument is modelled as a single temperature blackbody, \( M_{\text{inst}} = B(\nu, T_{\text{inst}}) \), where \( B(\nu, T) \) is the blackbody Planck function and \( T_{\text{inst}} \) is the temperature of the instrument enclosure in Kelvin (available from housekeeping telemetry). The instrument is usually at \( \sim 5 \) K and following Wien’s displacement law, the peak of the instrument emission is at \( \sim 600 \) \( \mu \)m, thus \( M_{\text{inst}} \) is much more significant for the longer wavelength SLW band than for the SSW band.

The telescope model used in the pipeline is a sum of two blackbody models, one for the primary and one for the secondary mirrors:

\[
M_{\text{tel}} = E_{\text{cont}}(t) \varepsilon_{1} (1 - \varepsilon_{2}) B(\nu, T_{M1}) + \varepsilon_{2} B(\nu, T_{M2}),
\]

(2)

where \( \varepsilon_{1} = \varepsilon_{2} \equiv \varepsilon(\nu) \) is the frequency-dependent telescope mirror emissivity, and \( T_{M1} \) and \( T_{M2} \) are the average temperatures of the primary and secondary mirrors, obtained via telemetry from several thermometers placed at various locations on the mirrors. The emissivity in equation (2) was measured for representative mirror samples pre-launch by Fischer, Klaasen & Hovenier (2004). For a dusty mirror \( \varepsilon \) is of the order of 0.2–0.3 per cent in the 200–600 \( \mu \)m band, with large systematic uncertainties. The only measured point in the SPIRE band, at 496 nm, has \( \varepsilon = 0.23^{+0.08}_{-0.12} \) per cent. Based on repeatability analysis of a number of ‘dark sky’ observations in Hopwood et al. (2014), the model was corrected by a small (sub 1 per cent) and mission-date dependent adjustment to the emissivity, \( E_{\text{cont}}(t) \).

During the Herschel mission around the second Lagrange point of the Earth–Sun system, the primary mirror temperature \( T_{M1} \) was of the order of 88 K and the secondary mirror \( T_{M2} \) was colder by 4–5 K, i.e. at around 84 K. Even with the low emissivity the telescope thermal emission is the dominant source of radiation recorded by the detectors. Only a few of the sky sources observed with the SPIRE spectrometer are brighter than \( M_{\text{tel}} \) : nearby large planets (Mars, Saturn) and the Galactic Centre.

The calibration of the FTS requires the derivation of \( R_{\text{tel}}, R_{\text{inst}}, M_{\text{tel}}, \) and \( M_{\text{inst}} \), as we can then recover the source intensity using

\[
I_{S} [\text{W m}^{-2}\text{Hz}^{-1}\text{sr}^{-1}] = \frac{(V_{\text{obs}} - R_{\text{inst}} M_{\text{inst}})}{R_{\text{tel}}}. \tag{3}
\]

Note that all quantities in equation (3) are frequency dependent and derived independently for each FTS band (see Fulton et al. 2014). As the two bands SSW and SLW overlap in 944–1018 GHz, the intensities in this region should match within the uncertainties.

The point-source calibration is built upon the extended-source calibration, using a suitable model of the emission of a point-like source. In the case of the SPIRE FTS, the primary calibrator is Uranus, which has an almost featureless spectrum in the FTS bands and a disc-averaged brightness temperature model known with uncertainties within \( \pm 3 \) per cent (ESA-4 model; Moreno 1998; Orton et al. 2014). The point-source conversion factor, \( C_{\text{point}} \), is derived as \( C_{\text{point}} = M_{U \text{uranus}}/\varepsilon_{\text{U \uranus}} \), where \( M_{U \text{uranus}} \) is the observed extended-source calibration intensity from the planet (following equation 3) and \( M_{U \text{uranus}} \) is the planet’s model. \( M_{U \text{uranus}} \) is converted from the disc-averaged brightness temperature model in units of K to units of Jy, using the planet’s solid angle, as seen from the Herschel telescope at a particular observing epoch (see H17 for details). Hence, \( C_{\text{point}} \) is in units of [Jy/(W m\(^{-2}\) Hz\(^{-1}\) sr\(^{-1}\)])

The point-source calibration was validated using Uranus and Neptune models, which showed an agreement within 3–5 per cent (Swinyard et al. 2014). Furthermore, the calibration accuracy was confirmed using a number of secondary calibrators (stars, asteroids) with the agreement at a level of 3–5 per cent between point-source calibrated spectra and the photometry from SPIRE point-source calibrated maps (Hopwood et al. 2015). Therefore, we consider the point-source calibration as well established and in this paper our focus is on the extended-source calibration.

3 CROSS-CALIBRATION WITH SPIRE PHOTOMETER

The SPIRE photometer and the FTS are calibrated independently and it is therefore important to cross-match measurements from observations of the same target. The cross-calibration can be considered as a critical validation of the different calibrations and whether their derived accuracies could be considered realistic. The cross-calibration in the case of point sources was already mentioned in the previous section, while in this section we restrict our discussion to the extended-source case.

The cross-calibration is performed between the extended-source calibrated spectra, obtained as described in Section 2, and the extended-source calibrated SPIRE photometer maps. These maps use detector timelines calibrated to the integrated signal of Neptune (Bendo et al. 2013) instead of the Neptune peak signal used for point-source calibrated maps. The arbitrary zero-level of each map is matched to the absolute zero level derived from Planck (Bertincourt et al. 2016). There is a good overlap of the SPIRE 350 \( \mu \)m band with the Planck-HFI 857 GHz band, and a relatively good overlap between the SPIRE 500 \( \mu \)m band and the Planck-HFI 545 GHz band. There is no Planck overlap for the SPIRE 250 \( \mu \)m band, so an extrapolation is used, based on a modified blackbody curve and the observed SPIRE 250 \( \mu \)m and Planck-HFI intensities (see H17 for more details). The overall uncertainty in the Planck-derived zero level is estimated at \( \sim 10 \) per cent, but for maps that are comparable in size to the Planck-HFI beam (FWHM \( \sim 5 \) arcmin, Planck Collaboration VII 2016) the uncertainty can be larger.

One of the most critical ingredients for extended-source calibration for any particular instrument is the knowledge of the beam and how the beam couples to a source (e.g. Ulich & Haas 1976; Wilson et al. 2013). Uncertainties on the beam solid angle or the beam profile as a function of frequency will lead to uncertainties in the derived quantities.

The SPIRE photometer beam maps were obtained using special observations of fine scans over Neptune and the same region of the sky at a different epoch when Neptune was no longer in the
field of view (i.e. the ‘shadow’ observation). Thanks to these two observations the photometer beams for the three bands have been characterized out to 700 arcsec and the beam solid angles are known down to the percentage level. Analysis of the beam maps for the three photometer beams indicates that the broad-band beams are unimodal and their cores are well modelled with 2D Gaussians (see H17; Schultz et al., in preparation).

On the other hand, the FTS beam was only measured out to a radial distance of 45 arcsec. The beam is multimoded and far from Gaussian, especially in the SLW band, which exhibits appreciable frequency-dependent beam FWHM variations (Makiwa et al. 2013). Hence, for sources with significant spatial brightness variation, the coupling with the beam is rather uncertain. Consequently, for the cross-calibration analysis, we need to identify spatially flat sources with as little source structure as possible within the FTS beam.

### 3.1 Selecting targets for cross-calibration

For all 1825 FTS observations performed with nominal bias mode (sparse and mapping modes, see H17), we extract an $11 \times 11$ pixel (66 arcsec × 66 arcsec) sub-image from the SPIRE 250 μm photometer map, centered on the SSW central detector coordinates. The SPIRE 250 μm beam FWHM is 18 arcsec and the largest SPIRE FTS beam has a FWHM of 42 arcsec (Makiwa et al. 2013), so the selected sub-image is bigger than the largest FTS beam FWHM for all frequencies. To characterize the surface brightness distribution in each sub-image we introduce the relative variation in the average level. Because of the Planck zero level normalization $I_{250} \gg 0$, no zero division effects are expected. To estimate the source flatness, we extract the central row and column from the sub-image and calculate two arrays of ratios: North–South:East–West and North–West:South–East. While either ratio can alone identify a vertical or horizontal gradient, the two ratios in combination from the sub-image and calculate two arrays of ratios: North–South:East–West and North–West:South–East. While either ratio can alone identify a vertical or horizontal gradient, the two ratios in combination can identify the surface brightness variation, the coupling with the beam is rather uncertain. Consequently, for the cross-calibration analysis, we need to identify spatially flat sources with as little source structure as possible within the FTS beam.

### 3.2 Synthetic photometry from extended-source calibrated spectra

To derive synthetic photometry from a spectrum we follow the approach explained in H17 and in Griffin et al. (2013). The total RSRF-weighted in-beam flux density from a source with spatial energy distribution $I_s(\nu)$ is

$$\bar{S}_S [\lambda] = \frac{\int_{\text{passband}} I_s(\nu) \eta(\nu) R(\nu) \Omega(\nu) \, d\nu}{\int_{\text{passband}} \eta(\nu) R(\nu) \, d\nu},$$

Here, $R(\nu)$ and $\eta(\nu)$ are the photometer spectral response function and the aperture efficiency for the passband. $\Omega(\nu)$ is the beam solid angle modelled with

$$\Omega(\nu) = \Omega(\nu_0) \left(\frac{\nu}{\nu_0}\right)^{2\gamma},$$

where $\Omega(\nu_0)$ is the beam solid angle derived from Neptune and $\gamma = -0.85$, $\nu_0$ is the adopted passband central frequency. The Neptune derived beam solid angles at the band centres (250, 350, 500) μm are $\Omega(\nu_0) = (469.35, 831.27, 1804.31)$ arcsec$^2$ (see H17).

A common convention in astronomy is to provide monochromatic flux densities or intensities at a particular central frequency $\nu_0$, assuming a source with a power-law spectral shape: $I(\nu) \propto \nu^{-\gamma}$. This convention is also used to calibrate the SPIRE photometer timelines. Hence, to convert $\bar{S}_S$ to monochromatic intensity $I_S(\nu_0)$ in [MJy sr$^{-1}$] for a source with $I(\nu) \propto \nu^{-\gamma}$ we use

$$I_S(\nu_0) = \text{KMonE}(\nu_0) \times \bar{S}_S,$$

where the conversion factors KMonE($\nu_0$) is

$$\text{KMonE}(\nu_0) = \frac{v_0^{-\gamma} \int_{\text{passband}} \eta(\nu) R(\nu) \Omega(\nu) \, d\nu}{\int_{\text{passband}} \eta(\nu) R(\nu) \, d\nu},$$

and the corresponding values are (91.567, 51.665, 23.711) in units of [MJy sr$^{-1}$ per Jy/beam] for the three photometer bands at (250, 350, 500) μm.

We use equations (4) and (6) to derive the synthetic photometry of extended-source calibrated spectra $I_s(\nu)$ from the two co-aligned central detectors of the two FTS bands. The error on the synthetic photometry is calculated by substituting $I_s(\nu)$ in equation (4) with $I_s(\nu) \pm \Delta I_s(\nu)$, where $\Delta I_s(\nu)$ is the standard error after averaging the different spectral scans in the pipeline (see Fulton et al. 2014 for details).\(^5\)

The 250 and 500 μm photometer bands are fully covered by the SSW and SLW spectra; however, the 350 μm band is mostly in SLW but a small fraction falls within SSW (see Fig. 2). For a source with $I(\nu) \propto \nu^{-\gamma}$, the underestimation of the synthetic photometry is ~1 per cent and for a $\nu^2$ spectrum it is overestimated by ~2 per cent. These are within the overall calibration uncertainties and consequently we do not stitch together the SSW and SLW spectra before deriving the synthetic photometry at 350 μm.

### 3.3 Comparison with the photometer

For each of the 24 flat sources we derive synthetic photometry as described in Section 3.2. The resulting values can be directly

\(^{3}\) Very few FTS observations have no associated SPIRE photometer map.

\(^{4}\) We do not include low resolution observations as in some cases the calibration introduces significant artefacts, mostly in the SLW band (Marchili et al. 2017).

\(^{5}\) This framework is implemented in the Herschel Interactive Processing Environment (HIPE) as a task spireSynthPhotometry(). The output of the task is the synthetic surface brightness values at 250, 350, and 500 μm in MJy sr$^{-1}$, for a monochromatic fully extended source with $R(\nu) \propto \nu^{-1}$.
Table 1. List of the final selection of spatially flat sources. The target name is that provided by the proposer. The equatorial coordinates RA and Dec. are for the central detector from the SSW array. For mapping we only used one FTS sparse snapshot out of 4 or 16 that were used to build the spectral cube. Only one SPIRE photometer and one PACS photometer OBSID are provided, although there can be multiple overlapping observations. If PACS and SPIRE photometer OBSIDs are the same then the observation was taken in Parallel Mode (see H17).

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Figure 1. (Left) Spatially flat extended source, rcw82off2 (see Table 1), with grey-scale image corresponding to the PACS 70 µm map and the 1 arcmin radius unvignetted FTS field of view shown as a red circle. The centre of the FTS field is marked with a ‘+’ sign. Note that the region appears as dark due to the very bright nearby rcw82; the peak surface brightness within the FTS footprint at 250 µm is more than 400 MJy sr⁻¹. (Right) Cas A – a supernova remnant shown compared to the corresponding extended-source calibrated photometer maps, by using a suitable aperture to take the average surface brightness. We use a square box aperture of 30 arcsec, which differs from the one used for the selection of extended and flat sources (Section 3.1). However, since we are averaging the surface brightness of flat extended sources then the choice of aperture is not important, as long as the size is comparable with the FTS beam.

Fig. 2 shows the extended-source calibrated spectrum produced with version 13.1 of the FTS pipeline for one of the flat sources (rcw82off2, ID8 in Table 1) and the derived synthetic photometry compared with the average surface brightness on 70 µm PACS data that was rejected because of its complex morphology although Δgmax = 0.07 and σI = 0.07.

Version 13.1 of the pipeline is the last one before the correction described in this paper was implemented.
The far-field feedhorn efficiency

The results shown in Fig. 3 (as well as the example in Fig. 2) indicate that in order to match the spectra with the photometry from the extended-source calibrated maps we need to apply a correction. We consider the SPIRE photometer extended-source calibration more straightforward than that of the spectrometer: simple beam profile, unimodal Gaussian beam and the beam solid angle is known down to <1 per cent uncertainty, and is consequently much more representative and robust. Moreover, the photometer maps are cross-calibrated with Planck-HFI. Therefore, the correction should be applied to the SPIRE FTS extended-source calibrated spectra.

The derived ratios, shown in Fig. 3, are a good match to the far-field feedhorn efficiency curve, \( \eta_{ff} \). The correction, \( \eta_{ff} \) was introduced in empirical form in Wu et al. (2013), where it was linked with two other corrections: the diffraction loss predicted by the optics model, \( \eta_{diff} \) (Caldwell et al. 2000) and the correction efficiency \( \eta_{c} \), with \( \eta_{ff} = \eta_{c} / \eta_{diff} \). As discussed in Wu et al. (2013), for point-like sources \( \eta_{c} \approx 1 \), while for extended sources \( \eta_{c} \ll 1 \) with the difference attributed to a combination of diffraction losses \( \eta_{diff} \) and different response of the feedhorns and bolometers to a source filling the aperture and to that of a point source.

The far-field feedhorn efficiency \( \eta_{ff} \) was measured by Chattopadhyay et al. (2003) but only for the SLW band (the two laboratory measurements are shown as red circles in Fig. 3). The empirical \( \eta_{ff} \) from Wu et al. (2013) is 10 per cent lower for SSW (shown as a dashed line in Fig. 3) with respect to the measured ratio at 250 \( \mu \)m. This 10 per cent is within the uncertainty of the 250 \( \mu \)m average ratio, however, the original empirical \( \eta_{ff} \) would introduce a significant discontinuity in the overlap region of the two FTS bands (944–1018 GHz). In order to avoid this inconsistency, \( \eta_{ff} \) was rescaled by 10 per cent for SSW, so that it matches the 250 \( \mu \)m ratio and also avoids the discontinuity. It is irrelevant to attribute this 10 per cent offset to any parameter in the optical model (\( \eta_{diff} \); Caldwell et al. 2000). The most likely interpretation is that some unknown effects in the complicated feedhorn-coupled system lead to a different response for fully extended sources only for SSW, which leads to \( \eta_{c} \approx 1.1 \) for SSW, while for SLW \( \eta_{c} = 1 \).

In practice, due to implementation considerations, we use the following empirical approximation based on the \( \eta_{ff} \) curves shown in Fig. 4 in Wu et al. (2013), with SSW rescaled by 10 per cent:

\[
\text{SLW} : \frac{1}{\eta_{ff}} = \frac{2.172}{v} - 1.47 \times 10^{-3}v, \\
\text{SSW} : \frac{1}{\eta_{ff}} = \frac{1.0857}{v} + 2.737 \times 10^{-4}v, \\
\]  

where \( v \) is the frequency in GHz. The two curves are shown in Fig. 3. And the corrected intensities are

\[
I'_{\text{ext}} = I_{\text{ext}} / \eta_{ff},
\]  

where \( I_{\text{ext}} \) is the extended-source calibrated spectrum from Swinyard et al. (2014) calibration [see also equation (3)]. Performing the same comparison for \( I'_{\text{ext}} \) with the extended-calibrated maps from the photometer for the 24 flat sources, we obtain the ratios as shown in Fig. 4. On average we see a good agreement at a level of 2–4
where $\Omega_{\text{beam}}(v)$ is the main beam solid angle.

Equation (11) should be valid for any instrument. And it is indeed the case for the SPIRE photometer, where the conversion from point-source to extended-source calibrated maps can be achieved by multiplication with $K_{\text{eff}}(v) = \Omega_{\text{pip}}/(\pi \Omega_{\text{beam}})$, where $\Omega_{\text{pip}}$ is the beam solid angle used in the data processing pipeline (see H17 for more details). The gain and aperture corrections already incorporated in the point-source calibrated timelines in the data processing pipeline.

The validity of equation (11) for the corrected extended-source calibrated spectra is demonstrated in Fig. 5 for a point source (Neptune) and an extended source from the sample of 24 spatially flat sources. In this case, the efficiency factor $\eta$ is actually the diffraction loss correction, $\eta_{\text{diff}}$ as derived by Caldwell et al. (2000), using a simple optics model, incorporating the telescope secondary mirror and mirrors support structures. For a point source on axis $\eta_{\text{diff}}$ is of the order of 75 per cent. We see that equation (11) is fulfilled at a level of ±5 per cent, if we exclude noisier regions close to the band edges (Fig. 5, bottom panels).

The noise that appears in the point-source converted spectra in Fig. 5 (cyan curves) reflects the small-scale characteristics of $R_{\text{tel}}$ that are inherently present in $I_{\text{rel}}$. The original point-source calibrated spectrum of Neptune (Fig. 5, left) has much less noise because the point-source calibration is based on the smooth featureless model spectrum of Uranus and consequently $C_{\text{point}}$ accounts for those small-scale features of $R_{\text{rel}}$. Therefore, the pipeline-provided point-source calibrated spectra are better products and they should be used, rather than converting the extended-source calibration with equation (11).

Interestingly, the missing correction for the old calibration of the FTS extended-source spectra is obvious, if we construct the ratio of the left-hand and right-hand side of equation (11), i.e. $f_X = \eta_{\text{diff}} I_{\text{spec}}/I_{\text{rel}}$. This ratio should be 1 if equation (11) is valid, but as shown in Fig. 3, the grey curves, which are the derived $f_X$ for all 24 flat sources with the old calibration, match well with the empirical $\eta_{\text{ff}}$ instead.

### 5 PRACTICAL CONSIDERATIONS

All extended-source calibrated spectra, regardless of the observing mode and the spectral resolution, are corrected for the missing far-field feedhorn efficiency (equation 9). Using those for analysis of extended sources is straightforward: measuring lines and the continuum, with results in the corresponding units of W m$^{-2}$ sr$^{-1}$. A large fraction of the sources observed with the FTS, however, are neither point-like nor fully extended, we call them semi-extended sources. The framework for correcting the spectra for this class of targets is presented in Wu et al. (2013) and implemented in HIPE as an interactive tool – the SEMIEXTENDEDCORRECTOR (SECT). There are two possible ways to derive a correction for the source size (and/or a possible pointing offset): starting from an extended-source or from a point-source calibrated spectrum (see Wu et al. 2013, equation 14). The SECT implementation in HIPE follows the procedure starting from a point-source calibrated spectrum. As the point-source calibration is not affected by the far-field feedhorn efficiency correction, described in Section 3.4, so there should not be any changes in the SECT-corrected spectra.

In cases when there is a point source embedded in extended emission, then the background subtraction should be performed using the point-source calibrated spectra, regardless of the fact that...
the background may be fully extended in the beam. If you perform the background subtraction using $I_{\text{ext}}'$, then you cannot any longer use $C_{\text{point}}$ to convert the background subtracted spectrum to a point-source calibrated one. Instead, you have to use equation (11), and as explained in Section 4, this will introduce unnecessary noise in the final spectrum.

The same consideration is applicable for semi-extended sources, where the first step before the correction should be the background subtraction and then proceeding with SECT, both steps should be performed on point-source calibrated spectra.

Careful assessment of the source extension is always necessary, because in some cases the source may fall in the extended-source category in continuum emission but semi-extended or point-like in a particular line transition. This will dictate which calibration to use and what corrections to apply to the line flux measurements.

Finally, if for some reason one needs to recover the spectrum with the original calibration following Swinyard et al. (2014), then $C_{\text{point}}$ and the point-source calibrated spectrum can be used: $I_{\text{ext}} = S_{\nu}/C_{\text{point}}$.

6 IMPLICATIONS FOR SPIRE FTS USERS AND ALREADY PUBLISHED RESULTS

The significant correction for the extended-source calibration scheme presented by this work, was implemented as of HIPE version 14.1, and has already been described in H17 since 2017 February. All analysis based on extended-source calibrated FTS spectra, produced prior to that version, will be affected by the significant and systematic shortfall of the old calibration. Any integrated line intensity or continuum measurements will be underestimated by a factor of 1.3–2 and using them to derive physical conditions in objects will be subject to corresponding systematic errors.

To illustrate the magnitude of the deviations on the derived physical characteristics with the old calibration, we performed a simple simulation using RADEX (van der Tak et al. 2007) for an emitting region, assuming $n(H_2) = 6.3 \times 10^3 \text{ cm}^{-3}$, column density of $10^{16} \text{ cm}^{-2}$, and kinetic temperatures of 100 K (green curve), 90 K (orange), and 80 K (blue). The 100 K SLED is multiplied by $\eta_{\text{ff}}$ and the new uncorrected SLED is shown with red points with error bars assuming a conservative 10 per cent uncertainty in line flux measurements.

If we observe a region with $T_{\text{kin}} = 100$ K, but use the old calibration, then the measured $^{12}$CO lines (the green curve) will

Figure 5. Left-hand panel: comparison of Neptune pipeline derived point-source calibrated flux density $S_{\nu}$ in Jy (thick black line) with the flux density derived from the extended-source calibrated intensity $S_{\nu}' = \eta_{\text{diff}} \times I_{\text{ext}}' \times \Omega_{\text{beam}}$ (cyan), i.e. equation (11). The relative ratio of $S_{\nu}'/S_{\nu}$ is shown in the bottom panel. The overall agreement, in the less noisy parts of the two bands, is within 5 per cent. Right-hand panel: the same comparison for a fully extended source.

Figure 6. $^{12}$CO spectral line energy distribution model from RADEX (van der Tak et al. 2007) for an emitting region, assuming $n(H_2) = 6.3 \times 10^3 \text{ cm}^{-3}$, column density of $10^{16} \text{ cm}^{-2}$, and kinetic temperatures of 100 K (green curve), 90 K (orange), and 80 K (blue). The 100 K SLED is multiplied by $\eta_{\text{ff}}$ and the new uncorrected SLED is shown with red points with error bars assuming a conservative 10 per cent uncertainty in line flux measurements.

Of 1.3–2 and using them to derive physical conditions in objects will be subject to corresponding systematic errors.

If we observe a region with $T_{\text{kin}} = 100$ K, but use the old calibration, then the measured $^{12}$CO lines (the green curve) will
be underestimated by a factor of $\eta_{ff}$; these are shown in Fig. 6 as red points with 10 per cent measurement errors. Obviously the red points do not match the RADEX models with $T_{\text{Kin}} = 100$ K, they are at least 2–3 $\sigma$ away from the correct input model for lines with upper J $\leq$ 8. While models with $T_{\text{Kin}}$ between 85 and 90 K are much closer to the ‘measurements’ and consequently the derived temperature from the red points will be significantly underestimated.

Using the old calibration for studies based on line-to-line or line-to-continuum measurement will not be significantly biased for SSW, because the variation of $\eta_{ff}$ with frequency within the band is small. However, the variation across SLW is significant and in this case using uncorrected data will lead to the incorrect results.

The $\eta_{ff}$ correction to extended-source calibrated spectra results in new values for the frequency-dependent additive continuum offsets and FTS sensitivity estimates (see Hopwood et al. 2015). The new offsets and sensitivities are presented in H17 and their tabulation is available in the Herschel legacy repository as Ancillary Data Products.  

The correction with $\eta_{ff}$ also introduces a new source of uncertainty to the overall calibration error budget for extended sources. The two measurement points for $\eta_{ff}$ in SLW band have errors of 3 per cent (Chattopadhyay et al. 2003), and we assume the same error is applicable for the SSW band. Therefore, the overall calibration accuracy budget for extended-source calibration will have to incorporate the 3 per cent statistical uncertainty on $\eta_{ff}$. As the correction is semi-empirical and based on cross-calibration with the SPIRE photometer, the more conservative estimate of the overall uncertainty is of the order of 10 per cent, to match the uncertainties on the derived photometry ratios (Fig. 4).

7 CONCLUSIONS

We introduce a correction to the SPIRE FTS calibration for the far-field feedhorn efficiency, $\eta_{ff}$. This brings the cross-calibration between extended-source calibrated data for the spectrometer and photometer in agreement at a 2–4 per cent level for fully extended and spatially flat sources. With this correction, the FTS point-source and extended-source calibration schemes are now self-consistent and can be linked together using the beam solid angle and a gain correction for the diffraction losses.

All SPIRE FTS extended-source calibrated products (spectra or spectral maps) in the Herschel Science Archive, processed with pipeline version 14.1 have already been corrected for $\eta_{ff}$. Spectra processed with earlier versions are significantly underestimated and consequently the results derived with the old calibration should be critically revisited. It is important to note that while the correction is close to a constant factor for the SSW band, this is not the case for SLW. Hence, even relative line-to-line or line-to-continuum analysis for SLW is affected.

We have not discussed any possible reason as to why the far-field feedhorn efficiency was not naturally incorporated in the extended-source calibration scheme. With Herschel no longer operational, it is not possible to take new measurements in order to check any hypothesis. We can only speculate about possible causes. One plausible reason is that the FTS beam, which was only measured out to a radial distance of 45 arcsec, compared to the 700 arcsec for the photometer, has an important fraction of the power distributed at larger distances, or in the side-lobes. Another possibility could be that the coupling of the two instruments to extended sources, viewed through the telescope, differs in an unknown manner such as small residual misalignment. Both these hypotheses could play a part in $\eta_{ff}$ not being naturally incorporated into then extended-source calibration. The bottom line, however, is that with this correction the FTS calibration is now self-consistent and the cross-calibration with the SPIRE photometer is in good agreement.

Ground-based measurements of lines or continuum, in frequency ranges that overlap with the large spectral coverage of the FTS, may provide further insights on the correctness of the extended-source calibration, although the direct comparison will not be straightforward due to the complications in observing very extended emission with ground-based telescopes.

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Most of the data processing and analysis in this paper was performed in the Herschel Interactive Processing Environment (HIPE, Ott 2010).

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8 See Appendix with a list of URLs for the data products.
APPENDIX: AVAILABLE DATA PRODUCTS

Many useful calibration tables are available in the Herschel Legacy Area at http://archives.esac.esa.int/hsa/legacy. Here, we only list those with relevance to the current paper.

(i) Planetary models:
Models for the primary calibrators (Uranus and Neptune) are available at http://archives.esac.esa.int/hsa/legacy/ADP/PlanetaryModels/

(ii) FTS sensitivity curves and additive continuum offsets:
The curves derived from the updated calibration are available at http://archives.esac.esa.int/hsa/legacy/ADP/SPIRE-SPIRE-S_sensitivity_offset/

(iii) Diffraction loss curves:
The correction $\eta_{\text{diff}}$ as presented in Wu et al. (2013), and based on the optics model from Caldwell et al. (2000) is available at http://archives.esac.esa.int/hsa/legacy/ADP/SPIRE/SPIRE_Diffraction_loss/

(iv) SPIRE photometer RSRFs:
The RSTFs $R(\nu)$ and the aperture efficiencies $\eta(\nu)$ are available at http://archives.esac.esa.int/hsa/legacy/ADP/SPIRE/SPIRE-P_filter_curves/

(v) SPIRE calibration tree:
The last one (SPIRE_CAL_14_3) as well as previous version of the calibration tables are available as Java archive files (jar) at http://archives.esac.esa.int/hsa/legacy/cal/SPIRE/user/